

16.2-3 The three first order operators of classical physics -

- The equations of classical physics (Fluid Dynamics, Electromagnetism etc.) are formulated in terms of three 1st order operators: (Grad, Div, Curl)

Defn: An operator is a function whose inputs & outputs are functional

$$(1) \nabla f = \left(\frac{\partial f}{\partial x}, \frac{\partial f}{\partial y}, \frac{\partial f}{\partial z} \right)$$

$$\nabla = \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right)$$

$$(2) \text{Div } \mathbf{F} = \frac{\partial M}{\partial x} + \frac{\partial N}{\partial y} + \frac{\partial P}{\partial z} = \nabla \cdot \mathbf{F}$$

$$(3) \text{Curl } \mathbf{F} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ M & N & P \end{vmatrix} = (N_y - N_z) \mathbf{i} - (P_x - M_z) \mathbf{j} + (M_y - P_y) \mathbf{k}$$

$$= \nabla \times \mathbf{F}$$

- Each operator ∇ , Div , Curl has a 3 dimensional version of the Fundamental Theorem of Calculus (FTC) that goes with it

②

- (1) $\nabla \Leftrightarrow$ "Conservation of Energy"
- (2) $\text{Div} \Leftrightarrow$ Divergence Theorem
- (3) $\text{Curl} \Leftrightarrow$ Stokes Theorem
- Recall: FTC says "the integral of a derivative reduces to an undifferentiated function evaluated on the boundary"

Simplest Case: Math 21B

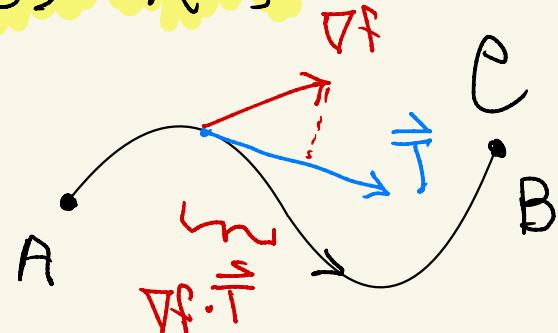
$$\int_a^b f'(x) dx = f(b) - f(a)$$

integral of the derivative on $[a, b]$

undifferentiated function evaluated on the boundary $x=a, b$

(1) ∇ : Conservation of Energy

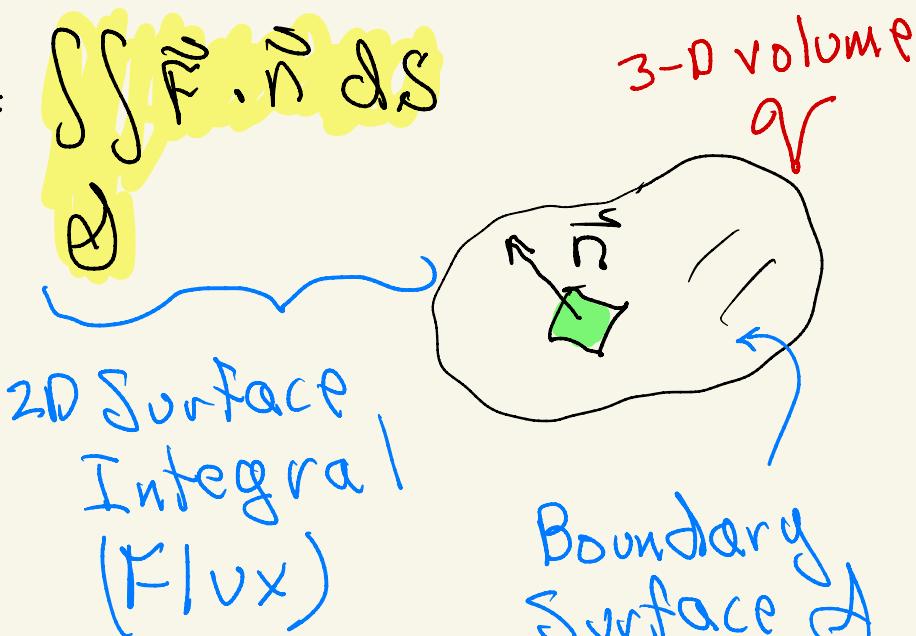
$$\int_C \nabla f \cdot \vec{T} \, ds = f(B) - f(A)$$



(2) Div : Divergence Thm

$$\iiint_V \operatorname{Div} \vec{F} \, dv = \iint_S \vec{F} \cdot \vec{n} \, ds$$

3-D triple integral
defined in §15



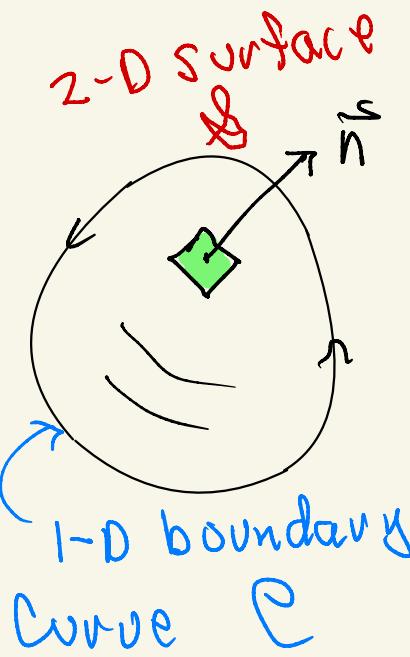
(3) Curl : Stokes Theorem

$$\iint_S \operatorname{Curl} \vec{F} \cdot \vec{n} \, ds$$

2-D surface
integral (Flux)

$$= \int_C \vec{F} \cdot \vec{T} \, ds$$

line integral
around closed
boundary



Our First Generalization of FTC (4)

(1) The FTC associated with the Gradient:

$$\nabla f = \left(\frac{\partial f}{\partial x}, \frac{\partial f}{\partial y}, \frac{\partial f}{\partial z} \right)$$

"input" ($f: \mathbb{R}^3 \rightarrow \mathbb{R}$) output $\nabla f: \mathbb{R}^3 \rightarrow \mathbb{R}^3$

We use the notation

$$\frac{\partial f}{\partial x} = f_x = \frac{\partial f}{\partial x}(x, y, z) = \lim_{\Delta x \rightarrow 0} \frac{f(x + \Delta x, y, z) - f(x, y, z)}{\Delta x}$$

We know (Math 21C) the Gradient takes a function $f: \mathbb{R}^3 \rightarrow \mathbb{R}$ (think of f as giving the temperature $f(\mathbf{x})$ at $\mathbf{x} = (x, y, z)$) and assigns to it the vector $\nabla f(x, y, z) = \overrightarrow{(M, N, P)}$ which points in direction of steepest increase of f .

- The first generalization of FTC involves the Gradient & we call it Conservation of Energy

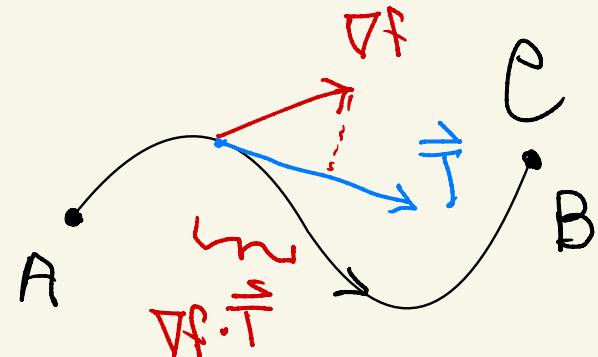
(1) ∇ : Conservation of Energy

$$\int_C \nabla f \cdot \vec{T} ds = f(B) - f(A)$$

Another way to say it:

We can evaluate the

line integral $\int_C \vec{F} \cdot \vec{T} ds$



if we can find an "anti-derivative" $\vec{F}(x)$

such that $\nabla f(x) = \vec{F}(x) = \overrightarrow{(M(x), N(x), P(x))}$

at every point $\vec{x} = \overrightarrow{(x, y, z)}$. In this case

$$\int_C \vec{F} \cdot \vec{T} ds = f(B) - f(A)$$

Defn: We say a vector field $\vec{F} = \overrightarrow{(M, N, P)}$ is conservative if there exists f st $\nabla f = \vec{F}$

Important: Most Vector Fields \vec{F} are NOT Conservative!

(6)

Questions to answer:

Q1: Why is it true?

Q2: Why is it called Conservation of Energy?

Q3: Given \vec{F} , how do you determine whether or not it is conservative?

It's easy to construct Examples:

Given $f(\vec{x})$, just calculate $\nabla \vec{F} = \left(\frac{\partial f}{\partial x}, \frac{\partial f}{\partial y}, \frac{\partial f}{\partial z} \right)$

Ex: Let $f(\vec{x}) = xy^2z^3 \Rightarrow \nabla \vec{F} = \left(y^2z^3, 2xyz^3, 3xy^2z^2 \right)$

Let C : straight line from $(0,0,0)$ to $(1,1,1)$.

Show $\int_C \vec{F} \cdot \vec{T} dS = f(1,1,1) - f(0,0,0) = 1 \cdot 1^2 \cdot 1^3 - 0 = 1$

ie Parameterize: $x = t, y = t, z = t \quad 0 \leq t \leq 1$
 $\vec{r}(t) = (t, t, t), \vec{v}(t) = (1, 1, 1)$

$\int_C \vec{F} \cdot \vec{T} dS = \int_0^1 \vec{F} \cdot \vec{v} dt = \int_0^1 (y^2z^3, 2xyz^3, 3xy^2z^2) \cdot (1, 1, 1) dt$

$\int_0^1 t^5 + 2t^5 + 3t^5 dt = \int_0^1 6t^5 dt = 6 \cdot \left[\frac{t^6}{6} \right]_0^1 = 1$

• Why is FTC-1 true?

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Ans: Chain Rule

$$\begin{aligned}\frac{d}{dt} f(x(t), y(t), z(t)) &= \frac{\partial f}{\partial x} \frac{dx}{dt} + \frac{\partial f}{\partial y} \frac{dy}{dt} + \frac{\partial f}{\partial z} \frac{dz}{dt} \\ &= \nabla f \cdot \left(\frac{dx}{dt}, \frac{dy}{dt}, \frac{dz}{dt} \right) \\ &= \nabla f \cdot \vec{v}(t)\end{aligned}$$

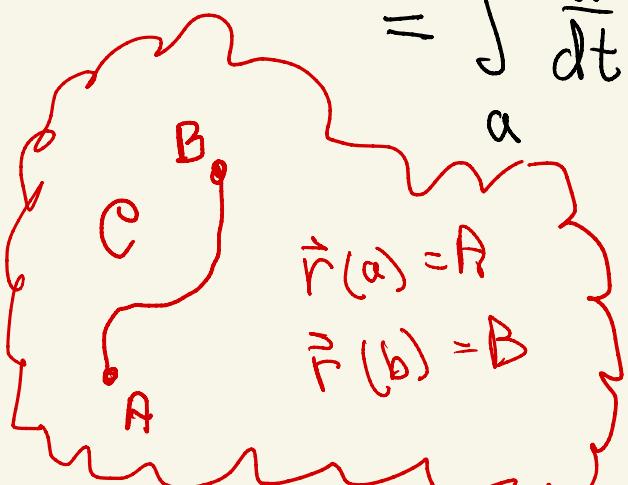
Therefore -

$$\begin{aligned}\int_C \vec{F} \cdot \vec{T} ds &= \int_a^b \vec{F} \cdot \vec{v} dt = \int_a^b \nabla f \cdot \vec{v} dt \\ \text{For } C: \vec{r}(t), a \leq t \leq b &\quad \text{assuming } \vec{F} \text{ conservative}\end{aligned}$$

$$= \int_a^b \frac{d}{dt} f(\vec{r}(t)) dt = f(\vec{r}(a)) - f(\vec{r}(b))$$

Math 21B
FTC

$$= f(B) - f(A)$$



✓

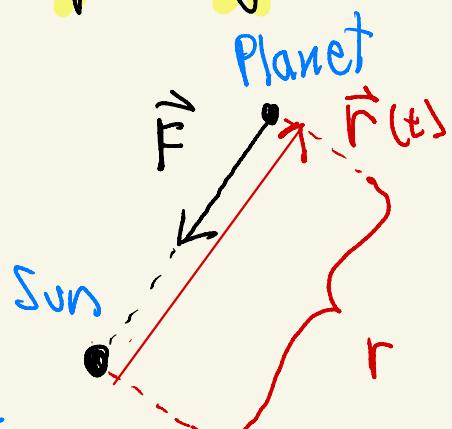
Q2: Why is it called Conservation of Energy? 8

Example: Newton's Theory of Gravity

Recall: Newton explained the motion of the planets by assuming the sun was pulling with an inverse square force:

$$\vec{F} = M_p \vec{a} = -G \frac{M_p M_s}{r^2} \frac{\vec{r}}{r}$$

magnitude in direction



Claim:

$$\nabla \frac{1}{|\vec{r}|} = -\frac{\vec{r}}{r^3}$$

$$\vec{r} = (x, y, z) = \vec{x}$$

$$|\vec{r}| = \sqrt{x^2 + y^2 + z^2} = r$$

Check: Let $f(x, y, z) = \frac{1}{|\vec{r}|} = \frac{1}{\sqrt{x^2 + y^2 + z^2}} = \frac{1}{|\vec{x}|}$

It suffices to show $\frac{\partial f}{\partial x} = -\frac{x}{|\vec{x}|^3}$

$$\text{I.e. Then } \nabla f = -\left(\frac{x}{|\vec{x}|^3}, \frac{y}{|\vec{x}|^3}, \frac{z}{|\vec{x}|^3}\right) = -\frac{1}{|\vec{x}|^3} \left(\frac{x}{|\vec{x}|}, \frac{y}{|\vec{x}|}, \frac{z}{|\vec{x}|}\right)$$

$$= -\frac{\vec{r}}{r^3}$$

(Cont.) We prove $\nabla \frac{1}{|x|} = -\frac{\vec{x}}{|x|^3}$ (9)

Calculations involving partial derivatives of $r = |x|$ are done so often it is very convenient to have a quick way to do them...

For this - note $\frac{\partial r}{\partial x} = \frac{x}{r} \rightarrow \frac{\partial r}{\partial y} = \frac{y}{r}, \frac{\partial r}{\partial z} = \frac{z}{r}$

i.e. $\frac{\partial r}{\partial x} = \frac{\partial \sqrt{x^2+y^2+z^2}}{\partial x} = \frac{1}{2} \frac{1}{\sqrt{x^2+y^2+z^2}} \cdot 2x = \frac{x}{r} \quad \checkmark$

Thus: $\frac{\partial}{\partial x} \frac{1}{r} = \frac{\partial}{\partial x} r^{-1}(x, y, z) = -1 \frac{1}{r^2} \frac{\partial r}{\partial x} = -\frac{1}{r^2} \frac{x}{r} = -\frac{x}{r^3}$

Therefore: $\frac{\partial}{\partial x} \frac{1}{r} = -\frac{x}{r^3}, \frac{\partial}{\partial y} \frac{1}{r} = -\frac{y}{r^3}, \frac{\partial}{\partial z} \frac{1}{r} = -\frac{z}{r^3}$
(Symmetry)

\Rightarrow (implies) $\boxed{\nabla \frac{1}{r^3} = -\frac{\vec{r}}{r^3}}$ \checkmark

Conclude:

$-\frac{GM_p M_s}{r^3} = GM_p M_s \nabla \frac{1}{r}$

Conclude = Newton's Gravitational Force Field is Conservative

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$$\vec{F} = -GM_s M_p \frac{\vec{r}}{r^3} = \nabla f = -\nabla P$$

$$f = GM_s M_p \frac{1}{r} \quad (f(x, y, z) = \frac{GM_s M_p}{\sqrt{x^2 + y^2 + z^2}})$$

Now we can apply FTC-1

$$\int_C \vec{F} \cdot \vec{T} ds = \int_C \nabla f \cdot \vec{T} ds = f(B) - f(A)$$

the work done

by Gravitation Force
as planet moves along C

the change in f
= minus the change in
"potential energy"

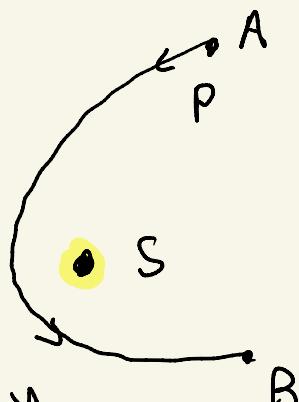
In physics: $P = \text{Potential Energy} = -f$
associated with a conservative force $\vec{F} = \nabla f$
is defined to be

$$P(x) = -f(x)$$

so $\vec{F} = -\nabla P$

\mathbb{C} minus sign

Picture : The change in potential energy $\Delta P = P(B) - P(A)$ keeps track of (i.e., is exactly equal to) minus the work done by \vec{F} along the motion.



Q: So why is the "work" defined by a line integral important in the first place?

Ans: If \vec{F} is the only force acting on the planet, i.e., then motion $\vec{r}(t)$ satisfies

$$M_p \cdot \ddot{\vec{r}}(t) = \vec{F},$$

the work done is also equal to the change in kinetic energy —

Theorem : If $\vec{F} = -GM_S M_p \frac{\vec{r}}{r^3}$, and $M_S \ddot{\vec{r}} = \vec{F}$,

then $\int_C \vec{F} \cdot \vec{ds} = \frac{1}{2} M_p V_B^2 - \frac{1}{2} M_p V_A^2 = \Delta KE$

the change in Kinetic Energy

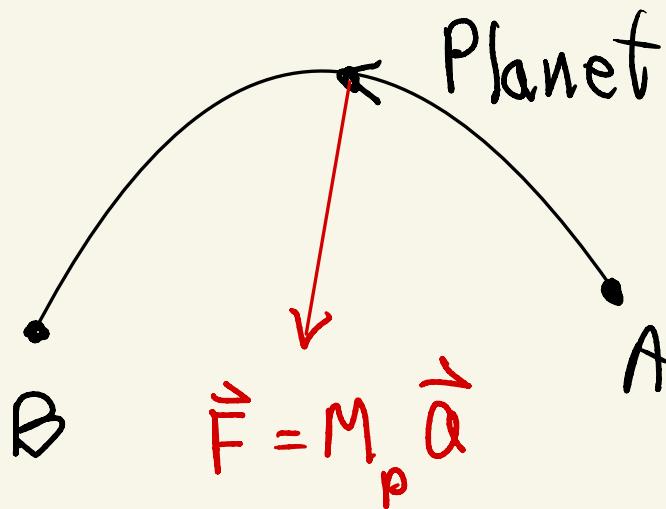
Proof of Theorem: Assume $M_p \vec{a} = \vec{F}$

(12)

where $\vec{F} = -GM_s M_p \frac{\vec{r}}{r^3}$. Then

$$\begin{aligned}
 \int_C \vec{F} \cdot \vec{T} \, ds &= \int_a^b \vec{F} \cdot \vec{v} \, dt = \int_a^b M_p \frac{d\vec{v}}{dt} \cdot \vec{v} \, dt \\
 &= M_p \int_a^b \frac{1}{2} \frac{d}{dt} (\vec{v} \cdot \vec{v}) \, dt = \frac{1}{2} M_p [\vec{v} \cdot \vec{v}] \Big|_{t=a}^{t=b} \\
 &= \frac{1}{2} M_p v_B^2 - \frac{1}{2} M_p v_A^2 = \Delta KE \quad \checkmark \\
 &\uparrow \qquad \qquad \qquad \uparrow \\
 &\vec{v} \cdot \vec{v} = v^2 \qquad \text{Change in Kinetic Energy}
 \end{aligned}$$

$$v(a) = v_A, v(b) = v_B$$



Conclude: If the planet moves according to

$$M_p \vec{a} = \vec{F}$$

where \vec{F} is the gravitational force of the sun

$$\vec{F} = -GM_s M_p \frac{\vec{r}}{r^3},$$

then along the motion

$$\Delta KE = \frac{1}{2} M_p V_B^2 - \frac{1}{2} M_p V_A^2 = \int_C \vec{F} \cdot \vec{T} ds = -(P(B) - P(A)) = -\Delta PE$$

OR: $\Delta KE + \Delta PE = \Delta \text{Energy} = 0$

We say energy is conserved all along the motion. This is Conservation of Energy

Conclude: FTC-1 $\int_C \vec{F} \cdot \vec{T} ds = f(B) - f(A)$

is the basis for the physical principle of Conservation of Energy -

② Summary - Newton unified all the laws of planetary motion known in his lifetime - namely Kepler's three laws - by postulating an inverse square force law between neighboring masses.

- $\mathbf{F} = M_p \vec{a}$ and $\vec{F} = -\frac{GM_p M_s}{r^2} \hat{r}$ led to

"three miracles", Kepler's three laws

- A "4th miracle" is that $\vec{F} = -\nabla P$

$$P(x) = -\frac{GM_p M_s}{r}$$

potential energy

So conservation of energy holds all along a planetary orbit $\vec{r}(t)$:

$$P(\vec{r}(t)) + \frac{1}{2} M \vec{v}(t)^2 = \text{constant}$$

This explains why orbits without a threshold energy remain trapped within the solar system